

which has less electron density if it is protonated.

The decay of the proposed radical cation VIII is an *intramolecular* electron-transfer process. The small negative activation entropy ( $-5.4 \text{ cal K}^{-1} \text{ mol}^{-1}$ ) can be regarded as further evidence that the large negative values of the entropy of activation which are observed for inner-sphere redox reactions of  $\text{Cr}(\text{H}_2\text{O})_6^{2+}$  with mononuclear<sup>14</sup> and binuclear<sup>15</sup> carboxylatocobalt(III)-ammine complexes are the result of concentrating the charge (2+, 2+ or 2+, 3+) in the dielectric medium as has been pointed out by Taube et al.<sup>16</sup>

The intrinsic reducibility of the aromatic, heterocyclic carboxylic acid appears to be the driving force for the rapid formation of relatively stable radical species, since they are generated when the unbound ligand as well as complexes A and B are reacted with  $\text{Cr}^{2+}$ . The stability of the Co(II-I)-Cr(III) radical cations II and VIII may at least in part be due to a low "electron permeability" of the carboxylate group.<sup>18</sup> This has been discussed previously for the relative stable radical cation formed by the reaction of  $\text{Co}^{\text{III}}(\text{NH}_3)_5(\text{O}_2\text{C}-\text{C}_6\text{H}_4-\text{NO}_2-p)^{2+}$  with the hydrated electron.<sup>17</sup>

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**Registry No.** Cr, 7440-47-3; pyrazinecarboxylic acid, 98-97-5; A, 37112-48-4; B, 62126-12-9;  $\text{Cr}(\text{H}_2\text{O})_6^{2+}$ , 20574-26-9; tri- $\mu$ -hydroxo-bis[triamminecobalt(III)] perchlorate, 22445-08-5.

### References and Notes

- (1) J. Halpern, *Q. Rev., Chem. Soc.*, **15**, 207 (1961).
- (2) A. Haim, *Acc. Chem. Res.*, **8**, 264 (1975).
- (3) H. Taube, "Electron Transfer Reactions of Complex Ions in Solution", Academic Press, New York, N.Y., 1970, p 66.
- (4) F. Nordmeyer and H. Taube, *J. Am. Chem. Soc.*, **90**, 1162 (1968).
- (5) E. S. Gould, *J. Am. Chem. Soc.*, **94**, 4360 (1972).
- (6) M. R. Hyde, K. L. Scott, K. Wieghardt, and A. G. Sykes, *J. Chem. Soc., Dalton Trans.*, 153 (1976).
- (7) E. S. Gould and H. Taube, *J. Am. Chem. Soc.*, **86**, 1318 (1964).
- (8) K. Wieghardt, *J. Chem. Soc., Dalton Trans.*, 2538 (1973); M. Hery and K. Wieghardt, *Inorg. Chem.*, **15**, 2315 (1976).
- (9) M. Linhard and H. Siebert, *Z. Anorg. Allg. Chem.*, **364**, 24 (1969).
- (10) E. S. Gould, *J. Am. Chem. Soc.*, **87**, 4730 (1965).
- (11) H. Spiecker and K. Wieghardt, in preparation.
- (12) H. Spiecker and K. Wieghardt, *Inorg. Chem.*, **15**, 909 (1976).
- (13) I. Bodek, G. Davies, and J. H. Ferguson, *Inorg. Chem.*, **14**, 1709 (1975).
- (14) M. B. Barrett, J. H. Swinehart, and H. Taube, *Inorg. Chem.*, **10**, 1985 (1971).
- (15) M. R. Hyde, K. Wieghardt, and A. G. Sykes, *J. Chem. Soc., Dalton Trans.*, 690 (1976).
- (16) H. Fischer, G. M. Tom, and H. Taube, *J. Am. Chem. Soc.*, **98**, 5512 (1976).
- (17) M. Z. Hoffmann and M. Simic, *J. Am. Chem. Soc.*, **94**, 1757 (1972).
- (18) H. Cohen and D. Meyerstein, *J. Chem. Soc., Dalton Trans.*, 2477 (1975).
- (19) E. Deutsch and H. Taube, *Inorg. Chem.*, **7**, 1532 (1968).
- (20) E. S. Gould, personal communication.

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## Crystal Circular Dichroism Spectra of the Tris(ethylenediamine)cobalt(III) Chromophore. Direct Measurement of Transitions to Both E and A<sub>2</sub> Excited Levels

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Crystal circular dichroism spectra of the  $\Lambda(+)_D$ -tris(ethylenediamine)cobalt(III) chromophore diluted in a host crystal of racemic  $2[\text{Ir}(\text{en})_3\text{Cl}_3 \cdot \text{NaCl} \cdot 6\text{H}_2\text{O}]$  were measured with light propagating both parallel and perpendicular to the threefold axis of the chromophore. The validity of the theoretical prediction that the trigonal components of the  $A_{1g} \rightarrow T_{1g}$  octahedral transition ( $A_1 \rightarrow E$  and  $A_1 \rightarrow A_2$ ) contribute to the rotatory strength with approximately the same absolute value, but opposite signs, is demonstrated experimentally, thus supporting the interpretation of the solution spectra as being residual wing absorptions arising from extensive cancelation of the rotatory strengths of these two transitions.

### Introduction

Optically active tris(diamine)cobalt(III) complexes usually exhibit two circular dichroism (CD) bands of opposite signs in the absorption region around  $21\,000 \text{ cm}^{-1}$ . These bands are normally considered to correspond to transitions to the  $E(D_3)$  and  $A_2(D_3)$  excited levels of  $T_{1g}(O_h)$  parentage in a  $D_3$  environment, although there has been some disagreement in the literature about these assignments. McCaffery and Mason,<sup>4</sup> Schäffer,<sup>5</sup> and Kuroda and Saito<sup>6</sup> consider the solution spectra of the  $\Lambda(+)_D$ - $[\text{Co}(\text{en})_3]^{3+}$  chromophore to correspond to residual wing absorptions resulting from extensive cancelation of the rotatory strengths of the two transitions, since the trigonal splitting is small. This argument is based on the assumption that solid-state and solution CD have the same physical cause, i.e., derive from a molecular property which is not greatly affected by the environment of the complex ion. Support for this assumption was provided several years ago by McCaffery et al.,<sup>7</sup> who compared solution and microcrystalline (KBr disks) CD spectra. Woldbye<sup>8</sup> and Dingle and

Ballhausen,<sup>9</sup> however, have argued that in contrast to the situation in the crystal, the complex ion may exist in two conformers in solution, so that the negative component in the solution spectrum of  $\Lambda(+)_D$ - $[\text{Co}(\text{en})_3]^{3+}$  could be due to the presence of conformations of the chelate rings in solution which differ from those occurring in the crystal. However, the discussion given by Hawkins,<sup>10</sup> the results of Piper and Karipides,<sup>11</sup> Sudmeier et al.,<sup>12</sup> and Harnung et al.,<sup>13</sup> and the calculations of Stickland and Richardson<sup>14</sup> provide good grounds for regarding the solution CD spectra of  $[\text{Co}(1,2\text{-diamine})_3]^{3+}$  in general as being due to residual wing absorptions. It was demonstrated by these authors that there are only minor differences between the CD spectra of the  $1e_1$  and  $1e_2$  conformers which occur in solution and that  $1e_1$ , the crystal conformer, shows two components in the CD spectrum. This means that results obtained for the solid phase may be directly related to solution and vice versa.

One of the difficulties encountered in attempts to resolve the controversy outlined above has been the fact that until very

Table I. Unit Cell Dimensions<sup>a</sup>

No.	Compd	<i>a</i> , Å	<i>c</i> , Å	10 <sup>24</sup> <i>V</i> , cm <sup>3</sup>	<i>d</i> <sub>calcd.</sub> <sup>b</sup> , g/cm <sup>3</sup>	<i>d</i> <sub>obsd.</sub> <sup>b</sup> , g/cm <sup>3</sup>
I	Mixed crystal	11.55	15.60	1802	1.97 <sup>b</sup>	1.95
II	2[Ir(en) <sub>3</sub> ]Cl <sub>3</sub> ·NaCl·6H <sub>2</sub> O	11.58	15.62	1815	2.06 <sup>b</sup>	2.05
III	2[Co(en) <sub>3</sub> ]Cl <sub>3</sub> ·NaCl·6H <sub>2</sub> O <sup>c</sup>	11.42	15.66	1770	1.61 <sup>b</sup>	1.61
IV	2 $\Lambda$ (+) <sub>D</sub> -[Co(en) <sub>3</sub> ]Cl <sub>3</sub> ·NaCl·6H <sub>2</sub> O	11.43	1/2 × 16.10	1/2 × 1820	1.56	1.56 <sup>d</sup>
V	[Co(en) <sub>3</sub> ]Cl <sub>3</sub> ·3H <sub>2</sub> O	11.44 <sup>f</sup>	15.48 <sup>f</sup>	1756	1.51	1.50 <sup>e</sup>

<sup>a</sup> Determined from x-ray powder photographs taken at 25 °C with Cu K $\alpha$  radiation using a focusing camera of the Guinier type, calibrated with silicon. All crystals belong to the hexagonal crystal system and the indexing of all powder photographs is very similar (see Table II).

<sup>b</sup> The density is calculated on the basis of two 2[M(en)<sub>3</sub>]Cl<sub>3</sub>·NaCl·6H<sub>2</sub>O units, i.e., four metal atoms per unit cell. <sup>c</sup> Space group *P*3̄c1, *Z* = 2. <sup>d</sup> Space group *P*3̄, *Z* = 1. <sup>e</sup> Space group *P*3̄c1, *Z* = 4. <sup>f</sup> Reference 23.

recently measurements of CD were restricted to solutions, glasses, and uniaxial crystals with light propagating along the optical axis, since the signal was otherwise distorted by interference from linear birefringence and linear dichroism. However, this problem was solved by Hofrichter and Schellman with the construction of a phase modulation spectrophotometer<sup>15</sup> and the technique was subsequently applied in measurements on an anisotropic crystal of an optically active copper(II) complex.<sup>16</sup> A discussion of the theoretical and practical aspects of the phase modulation technique in polarized spectroscopy is to be found in ref 16 and 17.

In the present work we provide evidence, based on measurements of CD parallel and perpendicular to the *c* axis of hexagonal mixed crystals of 2 $\Lambda$ (+)<sub>D</sub>-[Co(en)<sub>3</sub>]Cl<sub>3</sub>·NaCl·6H<sub>2</sub>O diluted in racemic 2[Ir(en)<sub>3</sub>]Cl<sub>3</sub>·NaCl·6H<sub>2</sub>O, to support the interpretation of the CD of microcrystals as being residual wing absorptions resulting from extensive cancelation of the rotatory strengths of the two first *D*<sub>3</sub> spin-allowed absorption bands of the  $\Lambda$ (+)<sub>D</sub>-[Co(en)<sub>3</sub>]<sup>3+</sup> chromophore.

### Experimental Section

The mixed crystals were prepared in the following way: 9.0 g (ca. 17 mmol) of [Ir(en)<sub>3</sub>]Cl<sub>3</sub>·aq, 3.0 g (51 mmol) of NaCl, and 1.0 g (1.17 mmol) of 2 $\Lambda$ (+)<sub>D</sub>-[Co(en)<sub>3</sub>]Cl<sub>3</sub>·NaCl·6H<sub>2</sub>O were dissolved in 30 mL of water by heating, and the solution was allowed to stand at room temperature for spontaneous crystallization. The best developed crystals were removed and wiped with tissue paper. The remaining crystals were dissolved by heating, and the solution was again allowed to stand for spontaneous crystallization. By repeating this procedure several times it was possible to harvest crystals developed preferentially with faces either parallel or perpendicular to the *c* axis of the hexagonal crystals. Investigation of a crystal with an electron microprobe showed that the surfaces are homogeneous, but because of deterioration of the crystal, the cobalt content could not be estimated reliably in this way. Atomic absorption spectroscopy of solutions of four crystals, precipitated at different stages, showed the crystals to contain 2.18, 2.14, 2.18, and 2.19% Co, respectively, corresponding to 19.8 mol % of the cobalt complex. Anal. Calcd for Ir<sub>1.6</sub>Co<sub>0.4</sub>NaCl<sub>12</sub>H<sub>60</sub>N<sub>12</sub>Cl<sub>7</sub>O<sub>6</sub>: C, 13.46; H, 5.65; N, 15.70; Cl, 23.17. Found: C, 13.55; H, 5.77; N, 15.88; Cl, 23.20; giving C:H:N:Cl = 12.00:60.9:12.06:6.96. A cobalt content in the mixed crystals of 23 mol % is calculated from the absorption spectrum of a solution of the latter. From the CD spectrum of the same solution a cobalt content of 22 mol % is calculated. We consider the atomic absorption measurements to be the most accurate and have therefore used a value of 20 mol % in calculating molar absorptivities. The CD spectrum of a solution of the mixed crystals exhibits a negative band with maximum at 268 nm, showing that the mixed crystals contain an excess of  $\Delta$ (+)-[Ir(en)<sub>3</sub>]<sup>3+</sup> (having<sup>18</sup>  $\Delta\epsilon_{\text{max}} = -0.9$  at 272 nm); i.e., some  $\Lambda$ -[Ir(en)<sub>3</sub>]<sup>3+</sup> in the racemic 2[Ir(en)<sub>3</sub>]Cl<sub>3</sub>·NaCl·6H<sub>2</sub>O lattice has been replaced in the mixed crystals by  $\Lambda$ -[Co(en)<sub>3</sub>]<sup>3+</sup>. This is supported by x-ray powder photographs of the compounds listed in Table I.

For compounds IV<sup>19,20</sup> and V<sup>21</sup> in Table I a complete x-ray structure analysis has been carried out by Nakatsu et al., who pointed out<sup>19</sup> that these structures are closely related; their close similarity is illustrated in Figure 1. Both unit cells have nearly the same *a*-axis dimensions and the length of the *c* axis for the optically active double salt is about half that for the racemic [Co(en)<sub>3</sub>]Cl<sub>3</sub>·3H<sub>2</sub>O; i.e., the basic structure of the optically active compound can be derived from

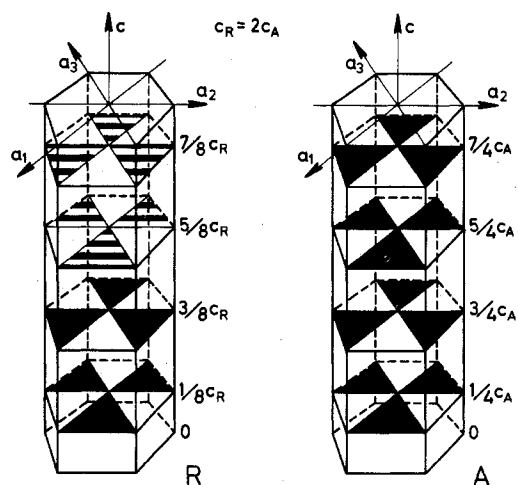


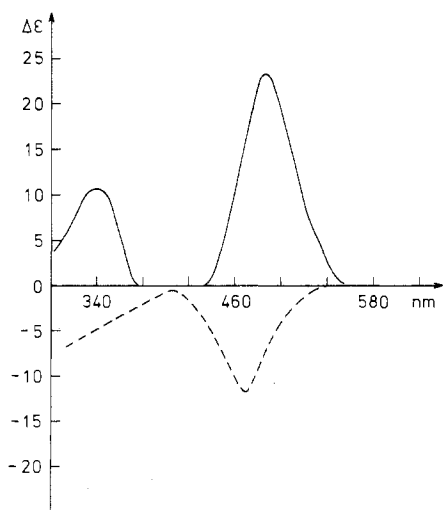
Figure 1. Resemblance of the structures of *rac*-[Co(en)<sub>3</sub>]Cl<sub>3</sub>·3H<sub>2</sub>O<sup>21</sup> (R) and 2 $\Lambda$ (+)<sub>D</sub>-[Co(en)<sub>3</sub>]Cl<sub>3</sub>·NaCl·6H<sub>2</sub>O<sup>19</sup> (A). The C<sub>3</sub> axes of the [Co(en)<sub>3</sub>(1el<sub>3</sub>)]<sup>3+</sup> units are parallel to the *c* axis of the crystal. The cobalt atoms of  $\Lambda$ -[Co(en)<sub>3</sub>]<sup>3+</sup> are situated at the center of the shaded triangles. The cobalt atoms of the  $\Delta$ -[Co(en)<sub>3</sub>]<sup>3+</sup> are situated at the center of the striped triangles.

Table II. Indexing of the Lines of the Powder Photographs of Compounds II and V of Table I

<i>hkl</i>	<i>I</i> <sub>obsd</sub>	2[Ir(en) <sub>3</sub> ]Cl <sub>3</sub> ·NaCl·6H <sub>2</sub> O		[Co(en) <sub>3</sub> ]Cl <sub>3</sub> ·3H <sub>2</sub> O		
		<i>d</i> <sub>obsd.</sub> <sup>a</sup> , Å	<i>d</i> <sub>calcd.</sub> <sup>a</sup> , Å	<i>I</i> <sub>obsd</sub>	<i>d</i> <sub>obsd.</sub> <sup>a</sup> , Å	<i>d</i> <sub>calcd.</sub> <sup>a</sup> , Å
10 $\bar{1}$ 0	vs	10.043	10.032	s	9.917	9.915
10 $\bar{1}$ 2	vs	6.171	6.162	s	6.100	6.099
11 $\bar{2}$ 0	vs	5.800	5.792	ms	5.736	5.724
11 $\bar{2}$ 1	vw	5.439	5.430			
20 $\bar{2}$ 0	w	5.019	5.016			
11 $\bar{2}$ 2	vw	4.656	4.652			
20 $\bar{2}$ 2	vs	4.221	4.220	vw	4.179	4.174
0004	s	3.902	3.904	m	3.872	3.868
11 $\bar{2}$ 3	vw	3.871	3.872	vw	3.832	3.832
21 $\bar{3}$ 0	vw	3.792	3.792			
21 $\bar{3}$ 1	w	3.684	3.685	w	3.642	3.642
10 $\bar{1}$ 4	s	3.637	3.638	mw	3.601	3.604
21 $\bar{3}$ 2	vs	3.411	3.411	m	3.373	3.373
30 $\bar{3}$ 0	vw	3.344	3.344			
11 $\bar{2}$ 4	ms	3.238	3.237			

that of the racemic compound by replacing the  $\Delta$  configurations by  $\Lambda$  configurations. From these observations and from the similarities of the relevant powder diagrams<sup>18</sup> we conclude that all of the compounds listed in Table I are closely related in structure. An example of this similarity is shown in Table II, and Table I summarizes the unit cell dimensions derived from the indexing of the powder diagrams for the compounds. Table I also shows the agreement between the densities observed for compounds I–III and those calculated assuming two formula units in the hexagonal cell.

In the above described manner we have thus achieved a solid-phase dilution of the cobalt complex which does not give rise to spectral interferences in the ligand field absorption region, since the iridium complex has very low absorption above 330 nm.<sup>18</sup> Crystals with



**Figure 2.** Experimental single-crystal circular dichroism spectra of  $2\Lambda(+)_D\text{-}[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot \text{NaCl} \cdot 6\text{H}_2\text{O}$  diluted in  $\text{rac-}2[\text{Ir}(\text{en})_3]\text{Cl}_3 \cdot \text{NaCl} \cdot 6\text{H}_2\text{O}$  measured parallel (—) and perpendicular (---) to the molecular  $C_3$  axis. Parallel maxima at 488 and  $\sim 335$  nm; perpendicular maximum at 470 nm.

developed  $\{0001\}$  and  $\{10\bar{1}0\}$ , i.e., the faces containing the  $a_1$ ,  $a_2$  and  $a$ ,  $c$  crystallographic axes, respectively, were polished down to an appropriate thickness (ca. 0.01 cm, giving rise to an absorbance of ca. 1 at 466 nm) using filter paper moistened with water. The crystals (two for each of the two orientations) were mounted on microscope slides and centered in the light beam of a microscope built into a phase modulation spectrophotometer as described elsewhere.<sup>15,16</sup> The results shown in Figure 2 are corrected for linear dichroism according to the expressions given in ref 16. The signs and the scaling of the spectra were established by the use of a pseudoisocyanin- $N,N'$ -diethyl iodide standard in a stretched PVA film,<sup>24</sup> measured both in a conventional dichrograph (Cary Model 6001 CD accessory) and in the phase modulation spectrophotometer.

The cobalt concentration (in mol/L) in the crystal is

$$c = \frac{Zn}{NV} \times 10^3 p \times 10^{-2}$$

i.e. the molar absorptivity (in L/(mol cm)) in the mixed crystal is given by

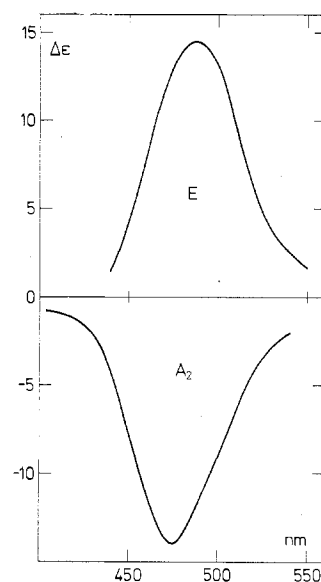
$$\Delta\epsilon = \frac{\Delta A}{cl} = \frac{\Delta ANV}{Zn(10)pl} \left( = \frac{\Delta AM}{nd(10)pl} \text{ cf. ref 9} \right)$$

where  $Z$  is the number of formula units per unit cell,  $n$  the number of chromophores per formula unit,  $N$  Avogadro's number,  $V$  the volume of the unit cell in  $\text{cm}^3$ ,  $p$  the cobalt content of the mixed crystal in mol %,  $\Delta A$  the measured circular dichroism,  $l$  the thickness of the crystal in cm,  $M$  the molecular weight of the mixed crystal, and  $d$  the density of the mixed crystal in  $\text{g}/\text{cm}^3$ . Rotatory strengths are evaluated from the expression<sup>25</sup>

$$R = 2.45 \times 10^{-39} \Delta\epsilon_{\text{max}} \frac{\Delta\nu_{1/2}}{\nu_0} \quad (\text{cgsu})$$

### Discussion

The crystal CD spectra of  $\Lambda(+)_D\text{-}[\text{Co}(\text{en})_3]^{3+}$  measured parallel and perpendicular to the molecular  $C_3$  axis, i.e., with light propagating into  $\{0001\}$  and  $\{10\bar{1}0\}$ , the faces containing the  $a_1$ ,  $a_2$  and  $a$ ,  $c$  crystallographic axes, respectively, are shown in Figure 2. The perpendicular spectrum is corrected for linear dichroism,  $A_c - A_a$ , which was measurable directly since the orientation of the crystal was fixed during the measurement so that the  $c$  axis of the crystal was parallel to the polarization direction of the light emerging from the linear polarizer.<sup>15,16</sup> Even though linear dichroism contributes,<sup>26</sup> the design of the phase modulation spectrophotometer makes this correction negligible.



**Figure 3.** Analysis of the  $21\,000\text{-cm}^{-1}$  region of the experimental single-crystal spectra (Figure 2) to give the E and  $A_2$  components. The experimental spectra have been corrected by the factor  $2/3$  according to eq 6–8 to bring the single-crystal data to a random orientation scale (cf. ref 5 and 27).

The separation of the maxima of the two CD signals in the  $21\,000\text{-cm}^{-1}$  region is ca.  $785\text{ cm}^{-1}$ . However, this is not a direct measure of the trigonal splitting since the positive curve corresponds to 3 times the average rotatory strength of the two E components (eq 6) and the negative curve to 3 times the average of one of the E components and the  $A_2$  component (eq 7 and 8). An analysis of the spectra places the E transition below the  $A_2$  transition in energy by ca.  $500\text{ cm}^{-1}$  (Figure 3). Accepting the idea of residual wing absorptions arising from cancellation of rotatory strengths, the spectra of Figure 2 should be correlated with a rotatory strength of  $4.3 \times 10^{-40}$  cgsu for the chromophore in random orientation.<sup>6,27</sup> In solution the total rotatory strength is defined by  $R = R(E) + R(A_2)$  with  $R$  ideally equal to zero for purely cubic parentage.<sup>5</sup>

If we place a Cartesian coordinate system in the molecular framework with the  $z$  axis along the molecular  $C_3$  axis and the  $y$  axis along a  $C_2$  axis, then the expressions for the rotatory strengths observed for a single crystal,  $R(\text{cr})$ , are as follows: for light propagating along the molecular  $C_3$  axis (i.e., with the two electric vectors defining the circularly polarized light spanning the  $x$ ,  $y$  plane)

$$R_{xy}(\text{cr}) = 1/2 [R_x(\text{cr}) + R_y(\text{cr})] \quad (1)$$

for light propagating along the  $y$  axis (i.e., with the two electric vectors defining the circularly polarized light spanning the  $z$ ,  $x$  plane)

$$R_{zx}(\text{cr}) = 1/2 [R_z(\text{cr}) + R_x(\text{cr})] \quad (2)$$

and for light propagating along the  $x$  axis

$$R_{yz}(\text{cr}) = 1/2 [R_y(\text{cr}) + R_z(\text{cr})] \quad (3)$$

If the molecules are in random orientation, a rotatory strength,  $R_{\text{rand}}$ , will be observed where

$$\begin{aligned} R_{\text{rand}} &= 1/3 R_{xy}(\text{cr}) + 1/3 R_{zx}(\text{cr}) + 1/3 R_{yz}(\text{cr}) \\ &= 1/3 R_x(\text{cr}) + 1/3 R_y(\text{cr}) + 1/3 R_z(\text{cr}) \end{aligned} \quad (4)$$

the factor  $1/3$  taking into account the fact that in a solution or a crystal powder on average only one-third of the molecules contribute in a given direction whereas in the single crystal all of the molecules are oriented in parallel directions. As rotatory strengths, absorptivities, etc. are usually referred to

the situation where the molecules are randomly oriented, the rotatory strengths are redefined as

$$R(E_x) = 1/3R_x(\text{cr}); R(E_y) = 1/3R_y(\text{cr}); R(A_2) = 1/3R_z(\text{cr}) \quad (5)$$

where  $R(E_x) = R(E_y) = 1/2R(E)$ . Equations 1-4 now become

$$R_{xy}(\text{cr}) = 3/2[R(E_x) + R(E_y)] = 3/2R(E) \quad (6)$$

$$R_{zx}(\text{cr}) = 3/2[R(A_2) + R(E_x)] = 3/2[R(A_2) + 1/2R(E)] \quad (7)$$

$$R_{yz}(\text{cr}) = 3/2[R(E_y) + R(A_2)] = 3/2[1/2R(E) + R(A_2)] \quad (8)$$

$$R_{\text{rand}} = R(E_x) + R(E_y) + R(A_2) = R(E) + R(A_2) \quad (9)$$

The rotatory strengths of the two  $A_{1g} \rightarrow T_{1g}(O_h)$  split components in Figure 2 are found experimentally to be  $R_{xy}(\text{cr}) = 64 \times 10^{-40}$  and  $R_{zx}(\text{cr}) = 30 \times 10^{-40}$  cgsu. The value for  $R_{xy}$  should be compared with experimental values of 79, 94.4, and 76.4 ( $\times 10^{-40}$  cgsu) found by McCaffery and Mason,<sup>4</sup> by Kuroda and Saito,<sup>6</sup> and by Judkins and Royer,<sup>28</sup> respectively, with the values 68<sup>4</sup> and 89.9<sup>6</sup> ( $\times 10^{-40}$  cgsu) reported for  $\Delta(+)_D\text{-[Co(en)}_3\text{]Br}_3 \cdot \text{H}_2\text{O}$ , with 62.3<sup>28</sup> and 57.2<sup>6</sup> ( $\times 10^{-40}$  cgsu) reported for  $\Delta(+)_D\text{-[Co(S)(+pn)}_3\text{]Br}_3$ , and with the theoretical values 75.5 and 95.9 ( $\times 10^{-40}$  cgsu) obtained for the  $\Delta\text{-[Co(en)}_3\text{]}^{3+}$  chromophore by Evans et al.<sup>29</sup> and by Mason and Seal,<sup>27</sup> respectively. The theoretical values have been corrected according to eq 6. Inserting our experimental values into eq 6, 7, and 9, we find  $R(E) = 43$ ,  $R(A_2) = -41$ , and  $R_{\text{rand}} = 1$  ( $\times 10^{-40}$  cgsu). We therefore conclude that the present measurements positively and independently demonstrate that both the microcrystalline and solution CD of the  $A_{1g} \rightarrow T_{1g}(O_h)$  transition indeed stem from cancelation of E and  $A_2$  bands. The value for  $R_{\text{rand}}$  is, of course, very uncertain as it is determined as the difference between two relatively large numbers.

Another interesting feature of the spectra shown in Figure 2 is the considerable intensity of the peak at  $\sim 335$  nm in the parallel spectrum. McCaffery and Mason<sup>4</sup> and Kuroda and Saito<sup>6</sup> found this to be of very low intensity. The discrepancy may be attributable to the different experimental techniques used, one or both of which may give rise to artifacts in the ultraviolet region. However, Barron<sup>30</sup> has discussed mechanisms whereby electric quadrupole moments may contribute to the rotatory strength of the high-frequency E transition of the  $[\text{Co(en)}_3]^{3+}$  chromophore when the latter is situated in an oriented crystal. The reason for the fact that we find a band with 10 times the rotatory strength found previously may very well be sought in our rather special host medium, in which the iridium centers give rise to large quadrupole transition moments, since such moments increase roughly as  $n^4/Z_{\text{eff}}^2$ .

As indicated by Barron the  $A_1 \rightarrow E(D_3)$  ( $A_{1g} \rightarrow T_{2g}(O_h)$ ) transition may "borrow" quadrupole transition moment from other E transitions, especially those of the crystal charge-transfer type which may extend over large molecular distances.

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**Registry No.** 2[ $\text{Ir(en)}_3$ ]Cl<sub>3</sub>·NaCl, 62126-09-4; 2[ $\text{Co(en)}_3$ ]Cl<sub>3</sub>·NaCl, 14881-39-1; 2 $\Delta(+)_D\text{-[Co(en)}_3\text{]Cl}_3 \cdot \text{NaCl}$ , 62210-58-6; [ $\text{Co(en)}_3$ ]Cl<sub>3</sub>, 13408-73-6.

## References and Notes

- Technical University of Denmark.
- Visitor at Institute of Molecular Biology, University of Oregon, Eugene, Oreg. 97403.
- University of Copenhagen.
- A. J. McCaffery and S. F. Mason, *Mol. Phys.*, **6**, 359 (1963).
- C. E. Schäffer, *Proc. R. Soc. London, Ser. A*, **297**, 96 (1967).
- R. Kuroda and Y. Saito, *Bull. Chem. Soc. Jpn.*, **49**, 433 (1976).
- A. J. McCaffery, S. F. Mason, and B. J. Norman, *Chem. Commun.*, 661 (1966).
- F. Woldbye, *Rec. Chem. Prog.*, **24**, 197 (1963).
- R. Dingle and C. J. Ballhausen, *K. Dan. Vidensk. Selsk., Mat.-Fys. Medd.*, **35**, No.12 (1967).
- C. J. Hawkins, "Absolute Configuration of Metal Complexes", Wiley-Interscience, New York, N.Y., 1971.
- T. S. Piper and A. G. Karipides, *J. Am. Chem. Soc.*, **86**, 5039 (1964).
- J. L. Sudmeier, G. L. Blackmer, C. H. Bradley, and F. A. L. Anet, *J. Am. Chem. Soc.*, **94**, 757 (1972).
- S. E. Harnung, S. Kallesøe, A. M. Sargeson, and C. E. Schäffer, *Acta Chem. Scand., Ser. A*, **28**, 385 (1974).
- R. W. Strickland and F. S. Richardson, *Inorg. Chem.*, **12**, 1025 (1973).
- H. J. Hofrichter and J. A. Schellman, *Jerusalem Symp. Quantum Chem. Biochem.*, **5** (1973); *J. Opt. Soc. Am.*, in press.
- H. P. Jensen, *Acta Chem. Scand., Ser. A*, **30**, 137 (1976).
- T. C. Troxell, H. P. Jensen, and J. A. Schellman, to be submitted for publication in *J. Opt. Soc. Am.*
- P. Andersen and F. Galsbøl, to be submitted for publication.
- K. Nakatsu, M. Shiro, Y. Saito, and H. Kuroya, *Bull. Chem. Soc. Jpn.*, **30**, 158 (1957).
- Y. Saito, K. Nakatsu, M. Shiro, and H. Kuroya, *Bull. Chem. Soc. Jpn.*, **30**, 795 (1957).
- K. Nakatsu, Y. Saito, and H. Kuroya, *Bull. Chem. Soc. Jpn.*, **29**, 428 (1956).
- K. J. Watson, private communication.
- P. Andersen, F. Galsbøl, and S. E. Harnung, *Acta Chem. Scand.*, **23**, 3027 (1969).
- B. Nordén, *Acta Chem. Scand.*, **27**, 4021 (1973).
- F. Woldbye, *Tech. Inorg. Chem.*, **4**, 286 (1965).
- S. Yamada and R. Tsuchida, *Bull. Chem. Soc. Jpn.*, **33**, 98 (1960).
- S. F. Mason and R. H. Seal, *Mol. Phys.*, **31**, 755 (1976).
- R. R. Judkins and D. J. Royer, *Inorg. Chem.*, **13**, 945 (1974).
- R. S. Evans, A. F. Schreiner, and P. J. Hauser, *Inorg. Chem.*, **13**, 2185 (1974).
- L. D. Barron, *Mol. Phys.*, **21**, 241 (1971).